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Optimizing the Magnetic Dipole-Field Source for Magnetically Guided Cochlear-Implant Electrode-Array Insertions

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Magnetic guidance of cochlear-implant electrode arrays during insertion has been demonstrated *in vitro* to reduce insertion forces, which is believed to be correlated to a reduction in trauma. In those prior studies, the magnetic dipole-field source (MDS) was configured to travel on a path that would be coincident with the cochlea's modiolar axis, which was an unnecessary constraint that was useful to demonstrate feasibility. In this paper, we determine the optimal configuration (size and location) of a spherical-permanent-magnet MDS needed to accomplish guided insertions with a 100 mT field strength required at the cochlea, and we provide a methodology to perform such an optimization more generally. Based on computed-tomography scans of 30 human subjects, the MDS should be lateral-to and slightly anterior-to the cochlea with an approximate radius (mean and standard deviation across subjects) of 64 mm and 4.5 mm, respectively. We compare these results to the modiolar configuration and find that the volume of the MDS can be reduced by a factor of five with a 43% reduction in its radius by moving it to the optimal location. We conservatively estimate that the magnetic forces generated by the optimal configuration are two orders of magnitude below the threshold needed to puncture the basilar membrane. Although subject-specific optimal configurations are computed in this paper, a one-size-fits-all version with a radius of approximately 75 mm is more robust to registration error and likely more practical. Finally, we explain how to translate the results obtained to an electromagnetic MDS.

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1. Introduction

Cochlear implants are neural-prosthetic devices implanted into the cochlea to directly stimulate the auditory nerve, bypassing the hearing mechanics of the auditory system and restoring effective hearing to those with profound sensorineural hearing loss. During the surgery, the surgeon either drills a hole in the cochlea (known as

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a cochleostomy) or makes an incision into the roundwindow membrane to insert an array of electrodes embedded in silicone, commonly referred to as the "electrode array" (EA), into the scala tympani chamber of the cochlea.

The insertion of the EA is known to cause intracochlear damage if the insertion forces exceed the inherent strength of the tissue [1]. The most traumatic damage occurs if the EA perforates the basilar membrane and deviates into the adjacent scala. The usual site for such trauma is along the basal (first) turn, where contact with the lateral wall sometimes deflects the electrode-array tip (EAT) out-of-plane and into the basilar membrane [1–3]. This type of trauma is a strong predictor for permanent loss of residual hearing [4]. Evidence also suggests that hearing outcomes improve by avoiding this

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type of trauma [5]. Preservation of residual hearing and reduction of insertion trauma is now a strong priority in cochlear implantation.

We previously developed a method to magnetically guide the EAT through the scala tympani during EA insertion, using an external magnetic dipole-field source (MDS) to apply the necessary magnetic torque to the EAT to actively bend it away from the walls of the scala tympani (Fig. 1) [6]. Repeatable automated insertions have recently been conducted in at-scale scala-tympani phantoms, designed with simulated cochleostomy and round-window openings [7], using clinical lateral-wall type EAs with magnets embedded at their tips [8]. Insertion forces were reduced by as much as 50% compared to nonguided insertions, and at the first turn in the basal plane, where a high percentage of basilarmembrane perforations occur, the EAT was never in contact with the lateral wall.

The bending torque that is applied to the EAT is given by the equation

$$\boldsymbol{\tau} = \mathbf{m} \times \mathbf{B},\tag{1}$$

where τ is the torque in units {N·m}, **m** is the dipole moment of the magnet at the EAT in units $\{A \cdot m^2\}$, and **B** is the magnetic field vector at the EAT in units $\{T\}$ (Fig. 1(d)). As illustrated in Fig. 1(b), the path of the EAT is mostly constrained on a plane orthogonal to the modiolus, particularly for the first turn, with a gradual climb in the direction of the modiolus: we will refer to this as the basal plane. Therefore, to bend the tip around the critical first turn where the EAT would normally contact the lateral walls, the component of the torque vector parallel to the modiolar axis should be maximized so as to bend the EAT principally within this plane. In other words, we want to minimize any component of the torque that could bend the EAT out of the plane and into the basilar membrane. This is accomplished by generating a magnetic-field vector that rotates on the basal plane while leading the EAT's dipole moment by up to 90°. All of our prior work achieved this by configuring the trajectory of the MDS to rotate about and translate along the modiolar axis (Fig. 1(a)), which we will refer to as the modiolar configuration. However, there is an advantageous reason for the MDS to approach the patient along a different trajectory: the size of the MDS can be reduced by positioning it closer to the cochlea. Our goal in this paper is to generate the required bending torque on the EAT from any location or trajectory that offers a clinical advantage, without the previous constraints placed on the MDS.

It has been shown that a rotating magnet, positioned anywhere relative to a fixed point in space, can generate at this fixed point a rotating magnetic field vector on any desired plane, simply by rotating the magnet about a unique axis-of-rotation [9]. This concept is illustrated in Fig. 2, where the rotating magnet is shown in three



Fig. 1. (a) In our original concept, magnetically guided insertions are achieved using three controlled degrees-of-freedom: (1) insertion of an EA with a magnetic tip; (2) rotation of the MDS about the modiolar axis to adjust the angle of the applied magnetic field B at the EAT; and (3) translation of the MDS along the modiolar axis to adjust the strength of the magnetic field at the EAT. (b) Close-up view of the scala tympani with the basilar membrane labeled (and shaded in white) so as to show its location relative to the EA and the MDS. (c) At each step of the insertion, the MDS applies torque to the EAT, bending the tip away from the cochlear walls. To minimize attractive force on the EAT toward the MDS, the angle between the magnetic orientations is maintained at approximately 90°. The increasing size of the MDS indicates that it is advancing toward the cochlea with increasing torque generation on the EAT. (d) Close-up view of basal plane, illustrating torque generation. Scala-tympani images are generated from software provided to the public by Eaten-Peabody Laboratory (Boston, MA).

example configurations, each with a unique axis-ofrotation ($\hat{\Omega}$) to generate a magnetic field vector at the cochlea that rotates on the basal plane, about the modiolar axis. Thus, from any arbitrary position outside the patient's head, the rotating MDS can generate the necessary rotating magnetic field for guided insertions. Note that in all the configurations shown, the direction from the MDS's center to the cochlea does not change. This allows the magnet to rotate about a constant



Fig. 2. Three example configurations in which a sphericalpermanent-magnet MDS translates relative to the patient's cochlea along a straight trajectory while rotating about some fixed axis $\hat{\Omega}$. The trajectories are not aligned with the surface normals of the patient's head, in general. The required size of the MDS will be determined by how close it can get to the cochlea without collision, which is depicted by the dashed circles. Note that configuration 1 corresponds to the modiolar configuration of Fig. 1(a).

axis-of-rotation while translating toward or away from the patient's head along a single degree-of-freedom. Part of the surgical planning could require the MDS to be positioned to some minimum clearance from the surface of the head with the linear stage extended to its maximum range of travel. Then, the linear stage could be retracted along the planned translation direction to the beginning of its range of motion. Such an alignment procedure would preclude any potential collision with the patient.

In this paper, we explore the optimal configuration in which the smallest MDS needed to accomplish guided insertions is determined along with its location relative to the cochlea. Throughout the optimization process, we assume that the MDS is a spherical permanent magnet. Spherical magnets have the desirable property that the point-dipole model perfectly describes their magnetic field [10]. In addition, our group has already developed a robotic end-effector that comprises a spherical permanent magnet [11]. Alternatively, an electromagnetic source such as an Omnimagnet [12], whose magnetic field is accurately described by the point-dipole model outside of its minimum bounding sphere, could be used in place of a spherical permanent magnet. In that case, it would not be necessary to translate the MDS relative to the patient's head to vary the field strength at the cochlea, because the strength of its magnetic dipole can be controlled directly. However, large currents will be necessary to generate the required dipole strength from an Omnimagnet of a given minimum bounding sphere compared to a spherical permanent magnet of the same size.

The outline of this paper is as follows. In Sec. 2, we present all of the necessary modeling equations that relate the size, strength, and location of the MDS to the magnetic field that can be generated at the cochlea. In Sec. 3, computed-tomography (CT) scans of 30 human subjects are segmented to create three-dimensional models of the surfaces of the subjects' heads, along with the location and orientation of the cochlea. In Sec. 4, we determine the optimal configuration (size and location) of the MDS for each of the 30 subjects. We compare the results to the MDS required in the modiolar configuration and find that the MDS size can be substantially reduced by moving it to the optimal location. Although we are primarily relying on magnetic torque applied to the EAT to bend the electrode array, in general there will also be a magnetic *force* applied to the EAT. In Sec. 5, we provide a conservative force estimate that indicates that magnetic forces are unlikely to apply any forces that might be dangerous to the delicate basilar membrane. In Sec. 6, we perform a sensitivity analysis to the size, strength, and positioning of the MDS. Although we compute subjectspecific optimal configurations in this paper, it is likely more practical to develop a one-size-fits-all MDS, which will be overdesigned for the majority of patients; in Sec. 7, we consider the robust placement of the MDS in light of this fact. In Sec. 8, we explain how to translate the optimization results to an Omnimagnet as the MDS. We conclude with a discussion of the practical application of our results in Sec. 9.

2. Magnetic Modeling

The necessary magnetic field to achieve successful guided electrode-array insertions has been determined by our prior work [8]. In that study, Flex-24 electrode arrays provided by MED-EL (Innsbruck, Austria) with a 4.73×10^{-5} A·m² permanent magnet embedded in the EAT were guided successfully through a plastic scalatympani phantom [7]. The maximum magnetic field used in those experiments was determined to be 80 mT and 100 mT if inserted through a simulated cochleostomy and round window, respectively. We will use 100 mT as the nominal magnetic-field requirement herein.

The magnitude of the dipole moment **M**, in units $\{A \cdot m^2\}$, for a spherical-permanent-magnet MDS of radius *r* in units $\{m\}$ is given by

$$\|\mathbf{M}\| = \frac{4\pi r^3 B_r}{3\mu_o},$$
 (2)

where $\mu_0 = 4\pi \times 10^{-7} \text{ T} \cdot \text{m} \cdot \text{A}^{-1}$ is the permeability of free space and $B_r = \mu_0 \tilde{M}$ is the residual flux density in units {T}, where \tilde{M} is the magnetization in units {A · m⁻¹}. B_r and \tilde{M} are intrinsic properties of the permanent magnet [10].

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The magnetic field vector \mathbf{B} in units {T} that is generated at the cochlea by a spherical magnet, modeled as a dipole \mathbf{M} , can be computed by the point-dipole equation

$$\mathbf{B} = \frac{\mu_0}{4\pi \|\mathbf{p}\|^3} \left(\frac{3(\mathbf{M} \cdot \mathbf{p})}{\|\mathbf{p}\|^2} \mathbf{p} - \mathbf{M} \right), \tag{3}$$

where **p** is the vector from the magnet's center to the cochlea in units $\{m\}$ [9]. We will define the location of the cochlea as the intersection of the modiolar axis and the basal plane [7].

We can simplify the use of (3) based on our intended use of the MDS. We intend to mount it on a one-degree-offreedom linear stage with its linear trajectory preplanned so as to avoid any contact with the patient at any point during its translation. During initial positioning, we envision the MDS to be placed as close to the surface of the head as allowable. This will represent the point on the trajectory where the distance between the MDS and the cochlea is shortest, as illustrated by the dashed circles touching the head surface in Fig. 2. In general, the shorter the distance between the MDS and the cochlea, the smaller the MDS can be and still generate the needed magnetic field. Now suppose that at its closest approach to the patient, the MDS rotates about the required rotation axis $\hat{\Omega}$, with M orthogonal to $\hat{\Omega}$, in order to generate a rotating field vector that lies on the basal plane. For every complete revolution of M about $\hat{\Omega}$, the field magnitude will vary based on the relationship between M and p, and it will reach a minimum (whenever $\mathbf{M} \cdot \mathbf{p} = 0$) of [9]:

$$\|\mathbf{B}\|_{\min} = \frac{\mu_0 \|\mathbf{M}\|}{4\pi \|\mathbf{p}\|^3}.$$
 (4)

We are searching for the location where the MDS can be smallest, so we cannot know the trajectory of the MDS in advance nor the relationship between **M** and **p**. In addition, if we insert the EA deep enough, the MDS will have to undergo a complete rotation. Therefore, to be conservative we will assume that the field magnitude is always at this minimum possible for a given **p**.

Assuming that the field magnitude necessary to achieve guided insertions of electrode arrays are known [8], then combining (2) and (4) yields the equation to compute the required minimum radius of the MDS.

$$r = \|\mathbf{p}\| \left(\frac{3\|\mathbf{B}\|_{\min}}{B_r}\right)^{1/3}.$$
 (5)

Taking each surface point on the head, we can simulate a spherical magnet touching the point and normal to the local surface as defined by the surface normal (see Fig. 3). By defining this surface point as **s** and its surface normal of unit length as $\hat{\mathbf{n}}$, the vector from the center of the magnet to the cochlea is

$$\mathbf{p} = -\mathbf{s} - (r+\delta)\hat{\mathbf{n}},\tag{6}$$

where *r* is the radius of the magnet and δ is the clearance from the surface point to the surface of the magnet along



Fig. 3. Vectors defining an arbitrary surface point **s** and its unit surface normal $\hat{\mathbf{n}}$ with respect to the cochlea. **p** is the vector from the center of the spherical magnet, which is touching the head at **s**, to the cochlea. θ is the angle between **s** and $\hat{\mathbf{n}}$.

 $\hat{\mathbf{n}}$. In the case where the magnet touches the patient, the clearance is $\delta = 0$.

The distance from the magnet center to the cochlea can be calculated using the (modified) law of cosines as

$$\|\mathbf{p}\|^{2} = \|\mathbf{s}\|^{2} + (r+\delta)^{2} + 2(r+\delta)s_{n},$$
(7)

where $s_n = \|\mathbf{s}\| \cos \theta$ and θ is illustrated in Fig. 3. s_n can also be considered the projection of $\|\mathbf{s}\|$ onto $\hat{\mathbf{n}}$.

Next, we rearrange (5) into the following form:

$$K \|\mathbf{p}\|^2 - r^2 = 0, \tag{8}$$

where the nondimensional term

$$K = \left(\frac{3\|\mathbf{B}\|_{\min}}{B_r}\right)^{\frac{2}{3}}.$$
(9)

expresses a ratio of the required magnetic field $(||\mathbf{B}||_{\min})$ to a magnetic property of the material (B_r) . By combining (7) with (8) and rearranging it into the form of a quadratic equation

$$(K-1)r^{2} + (2K\delta + 2Ks_{n})r + (K\|\mathbf{s}\|^{2} + K\delta^{2} + 2K\delta s_{n}) = 0$$
(10)

r required at some surface point **s** can be solved for using the quadratic formula. The standard quadratic formula yields two solutions, and the solutions are not guaranteed to be real. Fortunately, we can constrain some of the parameters, based on our application, to yield a valid solution to (10) (i.e. $r \in \mathbb{R}^+$): K > 0, $||\mathbf{s}|| > 0$, and $\delta \ge 0$.

We ignore all surface points with surface normals directed back toward the cochlea by selecting only the surface points in which $\mathbf{s} \cdot \hat{\mathbf{n}} \ge 0$ (i.e. $s_n \ge 0$). This typically eliminates some of the points on the ear or nose as shown in Fig. 4.

Finally, by further restricting 0 < K < 1, the solution to (10) always yields $r \in \mathbb{R}^+$ and is given by

$$r = \frac{K(\delta + s_n) + \sqrt{K(K(s_n^2 - \|\mathbf{s}\|^2) + \|\mathbf{s}\|^2 + \delta^2 + 2\delta s_n)}}{(1 - K)}.$$
(11)



Fig. 4. Surface rendering of a subject, with blue points indicating where $\mathbf{s} \cdot \hat{\mathbf{n}} < 0$ are excluded from the analysis.

The constraint on *K* implies that the required magnetic field strength at the cochlea must not be more than $\frac{1}{3}$ the residual flux density of the magnet.

$$3\|\mathbf{B}\|_{\min} < B_r. \tag{12}$$

Since the field measured on the surface of a spherical magnet falls within the range $\|\mathbf{B}\| \in \left[\frac{B_r}{3}, \frac{2B_r}{3}\right]$, (12) represents a conservative estimate of the field that can be generated by a spherical magnet in an arbitrary direction $\hat{\mathbf{p}}$.

3. Segmentation of CT Scans

Using (11) to compute the magnet size required at a given location relative to the patient's head requires a set of surface points and their associated surface normals. To our knowledge, a representative head-surface model of sufficient fidelity has not been published in which the location of the cochlea is also identified. In practice, patient-specific head surfaces can be generated from standard, preoperative CT scans. In this section, we

describe a generalized image-processing pipeline that we employed; we also provide the specific MATLAB commands used in Table A.1 in the appendix. The process yields a three-dimensional map of surface points and their surface normals for our analysis.

CT scans of 28 anonymous subjects were obtained from the University of Utah. Standard CT temporal-bone imaging procedures were used in each case. In addition, two sets of cadaver scans were obtained from the National Library of Medicine's Visible Human Project [13]. In all sets, the pixel resolution is less than 1 mm, but the slice resolution varied significantly from set to set. Data describing the gender and age of the 30 subjects, as well as the resolution of the CT scans, is provided in Table A.2 in the appendix. The CT data is available as DICOM binary data and packaged as 12-bit grayscale image intensity values and a header with meta-data regarding the CT scan and the imaging protocol used. The meta-data is needed to convert from image frame to the CT frame. We use the standard rightanterior-superior (RAS) convention in which +x is directed to the subject's right side, +y is directed anteriorly, and +z is directed superiorly.

All grayscale images are converted to black-and-white where the threshold to determine if the pixel should be considered black or white is set automatically by Otsu's method. Each slice of each set is then segmented to extract the boundary of the skull. We show an example of automatically generated boundaries, in red, overlayed on the grayscale image in Fig. 5(a).

The skull boundary is used to build a slice-specific image mask to remove all features outside this mask (Fig. 5(b)). The mask is a logical array where all locations inside the mask are set to true and all locations outside the mask are set to false. This mask can be visualized by



Fig. 5. Illustration of the main segmentation steps.

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mapping false and true to black and white, respectively. To remove the background, all pixels of the grayscale image that are outside this mask will have their intensity values set to zero, yielding an image of just the head without the background.

Next, this image mask is eroded (i.e. made smaller) for the purpose of handling internal features of the head, such as the nasal cavity (Fig. 5(c)). This is necessary because these features could be interpreted as part of the head surface. To handle this, we set all the intensity values for all the pixels inside this eroded mask to the maximum intensity value of the image slice.

The result is an image where the background is removed and the interior features are washed out, with a thin border of grayscale pixels at the head surface (Fig. 5(d)). The resulting segmented images are then stacked to create a three-dimensional array of intensity values per pixel and slice. This is illustrated in Fig. 5(e) (but only showing every tenth slice). This array of intensity values can also be regarded as a scalar function of three parameters I = f(x, y, z), where x and y are the pixel coordinates, z is the slice coordinate, and I is the scalar intensity value at those coordinates. As a result, this function can be interpolated to find the coordinate (x, y, z), where the intensity map transitions away from the background intensity. Interpolation then yields the approximate coordinates of the head surface, and this can be accomplished through the following steps. First, a homogenous transform is used to convert the pixel and slice coordinates from the image frame to the global CT frame. Then, the scalar function is filtered to generate smoother surface normals. Next, we interpolate for the coordinates at the desired intensity value. We chose an intensity value of 5 as representative of the transition from the background to the head surface (the images are provided as 12-bit unsigned integers, where 0 is mapped to black and 4095 is mapped to white). Finally, the surface normals corresponding to the surface points interpolated are generated. The combination of surface points and surface normals define the final head surface, as illustrated in Fig. 5(f). This pipeline is adapted from tutorials provided by Mathworks (Natick, MA) [14] for producing three-dimensional surfaces from medical images.

Finally, the location of the cochlea (as defined in Sec. 2) is determined for both sides of each subject (Fig. 6). In one study [15], the average angle between the modiolar axis and the transverse plane of the skull was determined to be nearly zero $(1.8^{\circ} \pm 3^{\circ}, n = 10)$. So, for simplicity, we will assume that the modiolar axis lies on the slice plane, where the cochlea is located. We examined each CT set to ensure the cochlea could be visualized in at least three of the slices. All sets summarized in Table A.2 of the appendix satisfied this criteria. Next, we chose to examine the slice with the largest cochlear footprint. This may yield a potential error of 1 slice along



Fig. 6. The location of the cochlea and the modiolar axis.

the inferior-to-superior coordinate. A line is then drawn through the basal chamber of the cochlea. A second line is drawn to estimate the central spiral axis (which we considered to be the modiolar axis). The intersection of these two lines at the chosen slice is considered to be the location of the cochlea. This visual approximation may yield an error of 2 to 3 pixels within the transverse plane of the subject, and we estimate a maximum localization error of 5 mm, largely due to the slice thickness. This method is imprecise compared with the automated algorithms [16] that would be used during surgical planning. However, in our opinion, this resolution is sufficient for the analysis in this paper given that the results are more sensitive to other factors as will be shown in Sec. 6, in which we include this localization uncertainty by treating it as an MDS registration error. Ultimately, the full pose of the basal plane is important for our magnetic-guidance approach because we want to apply the magnetic field in this plane, but it is not as important for the optimization described herein since (11) is independent of $\hat{\Omega}$.

An alternative method for computing the surface normals was attempted on the boundary points directly (Fig. 5(a)) so as to avoid the steps after boundary segmentation, but we found the surface normals to be more noisy and settled on the method described above. However, the possibility of noisy surface normals motivated an alternative method, independent of surface normals, to verify the results computed by (11); this is described in Sec. 4.1.3.

4. Determining the Optimal Configuration

4.1. Methods

4.1.1. Using surface normals

For each surface point obtained from the segmentation, and that satisfy $\mathbf{s} \cdot \hat{\mathbf{n}} \ge 0$, the MDS radius *r* is computed by (11). To ensure that this MDS does not collide with any point on the head surface, the corresponding \mathbf{p} is

computed using (6), assuming $\delta = 0$. If $\|\mathbf{p} + \mathbf{s}\| \ge r$ is true over the entire set of all surface points \mathbf{s} , no collision occurs, and this configuration is valid. The valid surface point that yields the smallest r is considered optimal. The corresponding \mathbf{p} determines its optimal placement relative to the cochlea.

4.1.2. On the modiolus

For comparison, we also determine the smallest MDS allowable if it is constrained to be on the modiolar axis, as in our prior work. Points are sampled on the modiolar axis at increments of 0.1 mm and represent the set of hypothetical locations for the MDS. At each sampled point, the smallest MDS (of radius *r* with $\delta = 0$) needed to generate the required magnetic field is computed using (5), where $\|\mathbf{p}\|$ is simply the distance from the sampled point to the cochlea. The last step is to perform collision detection at each sampled point to determine if the MDS will collide with any of the surface points on the head. If there are no collisions, then the sampled point is valid. The same condition is applied for collision detection as in Sec. 4.1.1. The valid point yielding the smallest radius is the smallest MDS allowable for the modiolar configuration.

4.1.3. Using brute force

To verify the results computed by (11), an alternative method that is independent of surface normals is also conducted. The method is identical to that described in Sec. 4.1.2, except that the entire space exterior to the head is uniformly sampled at the resolution of the subject-specific CT scan. Thus, for each sampled point, there will be an associated radius *r* that represents the smallest MDS required to generate $\|\mathbf{B}\|_{\min}$ at the cochlea if the MDS were centered at that sampled point. The optimal configuration using this method is the sampled point that yields the smallest *r* without any collisions.

4.2. Results

4.2.1. Using surface normals

The compiled results over the entire data set of subjects is presented in Table A.3 in the appendix. The optimal location is identified by the distance $||\mathbf{p}||$ and the direction from the cochlea to the MDS center. The direction can be expressed compactly by the unit vector $-\hat{\mathbf{p}}$; recall that $\hat{\mathbf{p}}$ is the unit vector in the opposite direction (i.e. from the MDS center to the cochlea). All unit vectors assume the RAS convention. Since the optimal MDS location is associated with a surface point, this is also provided as a distance ($||\mathbf{s}||$) and direction ($\hat{\mathbf{s}}$) from the cochlea. The angle between this surface point and its

surface normal (θ) is also listed. r and r_m are the computed radii at the optimal and modiolar configurations, respectively. At the bottom of Table A.3 are some basic statistics for the entire set.

Across all the subjects analyzed, there is a 19 mm range in the radii of the optimal MDS (55.0–73.5 mm). There is a 32 mm range in the distances from the cochlea to the MDS center (93.3–125 mm); this is the most variable of the parameters listed.

The optimal MDS placement is mostly lateral-to and slightly anterior-to the cochlea. This is a fortuitous result because the typical approach to the cochlea requires an incision behind the ear [17], making it impractical to position the MDS behind the ear. Our results indicate that the optimal configuration will not interfere with the surgical insertion. There does not seem to be a clear trend for the *z*-coordinate of the MDS as there seems to be a balance in the instances, where the MDS should be placed superior to or inferior to the cochlea. The direction to the optimal surface point tends to be more anterior than the direction to the MDS center.

4.2.2. On the modiolus

By moving the MDS away from the modiolar configuration, its size can be significantly reduced. For example, in the case of P28 shown in Fig. 7, the radius is reduced by half, resulting in an eight-fold decrease in the volume. On average, over our entire 30-subject data set, the radius is reduced by approximately 43% (see Table A.3), yielding an approximately five-fold decrease in the volume of the magnetic sphere on average. Also, the range and standard deviation in radius values are cut approximately in half in the optimal configuration. There is variability in the side-to-side optimal results within each individual, but in the aggregate the left and right values were very similar. If we consider only the maximum radii found over our entire 30-subject data set, as we might in the design of a one-size-fits-all system, we find that the radius computed in the modiolar configuration is 80% larger than it needs to be (132 mm versus 73.5 mm).

4.2.3. Using brute force

A table of results for the brute-force method, similar to Table A.3, is provided in Table A.4 in the appendix. With the exception of a few cases, the brute-force search method yielded slightly smaller radii values than the surface-normals method. On the aggregate, however, the surface-normals method yields an average radius of only about 1 mm larger, making it slightly conservative and thus better for our purpose. The main drawback to the brute-force method is that computation times are considerably longer because it typically yields approximately 1000 times more points to examine than the surfacenormals method.



Fig. 7. Three orthogonal views of a single subject showing how the size of the MDS is drastically reduced by moving it away from the modiolar axis (large MDS) to the optimal configuration (small MDS). Black lines are drawn from each of the MDS centers to the cochlea (which are all shown with black dots). The spatial location of the cochlea can be visualized as the intersection of these two black lines in each view (and can be seen directly in the upper-left view). A blue arrow indicates the surface normal at the optimal surface point (which is shown with a blue dot). These results are typical of all human subjects.

5. Magnetic Force

In all prior work, the MDS used to generate the magnetic field at the cochlea was assumed to be in the modiolar configuration. One benefit of this configuration is that the magnetic force that could potentially attract the EAT into the basilar membrane was negligible and could be disregarded during the insertion [6]. However, at the optimal configuration, this assumption should be verified since avoiding trauma to the basilar membrane is widely considered vital for hearing preservation [4]. The threshold for puncturing the basilar membrane has been measured to be approximately 42 mN [18]. We compute the worst case and assume that the entire force vector is directed into the basilar membrane. While this is not accurate, it does present a conservative upper bound.

The magnetic force applied to the tip is given by

$$\mathbf{f} = \nabla(\mathbf{B} \cdot \mathbf{m}) = \left[\frac{\partial}{\partial x} \mathbf{B} \ \frac{\partial}{\partial y} \mathbf{B} \ \frac{\partial}{\partial z} \mathbf{B}\right]^{\mathrm{T}} \mathbf{m}, \qquad (13)$$

where **B** is computed by (3) and **m** is the dipole moment of the EAT. For a given $||\mathbf{p}||$, the largest magnitude and spatial derivative of the field vector occurs along the dipole axis of **M** (i.e. where **M** and **p** are parallel). The maximum possible force magnitude at a given $||\mathbf{p}||$ is therefore computed by differentiating $\|B\|$ with respect to $\|p\|$, with M parallel to p, yielding

$$\|\mathbf{f}\| = \left\|\frac{\partial}{\partial \|\mathbf{p}\|} \left(\frac{\mu_0 \|\mathbf{M}\| \|\mathbf{m}\|}{2\pi \|\mathbf{p}\|^3}\right)\right\| = \frac{3\mu_0 \|\mathbf{M}\| \|\mathbf{m}\|}{2\pi \|\mathbf{p}\|^4}.$$
 (14)

As a nominal value for $||\mathbf{m}||$, we will use the magnets embedded in the EAT in our prior work [8] which has been determined to be $4.73 \times 10^{-5} \text{ A} \cdot \text{m}^2$. This represents the combined magnetic dipole for two 0.41-mm-long by 0.25-mm-diameter cylindrical magnets made of grade N52 NdFeB.

At the optimal position, for all subjects and sides examined, the maximum magnetic force possible never exceeded 0.31 mN, with an average and standard deviation of 0.267 mN and 0.019 mN, respectively. For comparison, we also compute the maximum magnetic force possible if the MDS were positioned at the modiolar configuration for each subject and found that the average and standard deviation is 0.150 mN and 0.012 mN, respectively. Note that by moving the MDS from the modiolar configuration to the optimal configuration, the maximum possible magnetic force averaged over the set is increased by nearly 80%. In context however, the force is still over 100 times smaller than what is needed to puncture the basilar membrane. Therefore, it seems reasonable to state that any additional magnetic force that may pull the EAT into the basilar membrane arising from the MDS being moved to the optimal configuration does not truly pose any appreciable risk to the basilar membrane.

6. Sensitivity Analysis

Achieving the optimal configuration requires perfect registration of the MDS with respect to the cochlea, and hardware implementations will have practical design constraints. A sensitivity analysis to registration, magnetic material properties, and any clearance requirements between the MDS and the patient would be useful. This is motivated by the reality that these factors may reduce the magnetic field generated at the cochlea and thereby limit the amount of insertion-force reduction that can be achieved with magnetically guided insertions.

For each parameter of interest, using (11), r can be solved for over the possible range of the parameter while fixing the remainder of the inputs to their nominal values. With this in mind, in this section, we first provide the nominal values used by the search method, as well as the range of expected values for those parameters. We conclude with plots that show the sensitivity of r to the input parameters within the range defined. These plots are provided to enable readers to get a quick sense of the how the optimal MDS can vary as a function of the various parameters. However, we do not recommend using these plots in a formal design; rather, the equations in Sec. 2 should be used directly.

The nominal value for B_r is 1.465 T, which corresponds to a NdFeB magnet of grade N52. This is the highest grade of magnetic material available, and it is what one would likely choose. The range of values for grade N35–N52 is 1.19–1.465 T [19]. We chose to use a nominal value of 100 mT for the required magnetic field strength $\|\mathbf{B}\|_{min}$. For the range of $\|\mathbf{B}\|_{min}$, we will explore 80–120 mT, representing a 20% variation in that parameter.

The nominal value for the gap (δ) was set to zero and describes the situation where the spherical magnet just touches the head. However, the spherical magnet itself will likely be contained in a housing. An existing MDS in our lab, using a 50-mm-diameter N42 NdFeB sphere, has been designed with a 7-mm-thick housing [11]. So, if this MDS were to be used, then $\delta = 7$ mm, assuming the housing of the MDS touches the head. We will use a range of values from 0–25 mm for δ to conduct the sensitivity analysis.

In our context, registration error can broadly be considered as the position error in the initial placement of the MDS center, relative to the cochlea. This error can be due to either misplacement of the MDS, or incorrectly estimating the location of the cochlea. We believe the contribution from errors in locating the cochlea will be negligible since Bell et al. [20] have already demonstrated successful automated insertions in temporal bones to sub-mm accuracy based on CT localization of the cochlea. For this reason, sensitivity to registration can be assessed by evaluating the effect on r due to error in $\|\mathbf{p}\|$ as indicated by (5). Although the optimal results found in Table A.3 were computed using (11), which utilized the surface normal, r is most sensitive to distance changes along $\hat{\mathbf{p}}$, which makes consideration of error along $\hat{\mathbf{p}}$ a conservative bound on registration error in any direction by the same amount. In Fig. 8, r is computed based on the change from the optimal $\|\mathbf{p}\|$ for each human subject (Table A.3) over the range of ± 10 mm.



Fig. 8. *r* as a function of $\Delta ||\mathbf{p}||$ for all subjects analyzed. C1 and C2 results are shown in red, with all others shown in black. The blue-dashed and blue-dotted lines indicate two and three standard deviations above the mean, respectively.



Fig. 9. (a) Expected range of *K* given the range of $||\mathbf{B}||$ considered, for commercially available permanent-magnet grades of the MDS. The blue vertical line is the nominal value for $||\mathbf{B}||$, and N52 is the nominal grade. The resulting range of *K* is used in (b), which shows *r* as a function of *K* for all subjects analyzed. C1 and C2 results are shown in red, with all others shown in black. The blue vertical line is the nominal value for *K*, which is the intersection of $||\mathbf{B}|| = 100 \text{ mT}$ and N52.

The results for each subject and each side are plotted together on a single graph, assuming the nominal value of $\|\mathbf{B}\|_{\min}$ and B_r .

Sensitivity results pertaining to the MDS hardware design are grouped into two categories. The first category represents the effect from the magnetic field desired and the magnetic property of the MDS, as expressed by the nondimensional parameter *K*. In Fig. 9(a), we display K over the range of $||\mathbf{B}||$ outlined earlier (i.e. 80-120 mT) and the range of magnetic grades commercially available. This range of *K* is then used to compute the values of *r* in Fig. 9(b). This is computed for each subject and each side separately and plotted together on a single graph. We assume the nominal value of δ , $||\mathbf{s}||$, and θ for these results.

The second category represents the effect from increasing the clearance between the magnet and the surface of the subject. In Fig. 10, *r* is computed as a function of δ over the range defined earlier (i.e. 0–25 mm). This is



Fig. 10. *r* as a function of δ for all subjects analyzed. C1 and C2 results are shown in red, with all others shown in black. The blue-dashed and blue-dotted lines indicate two and three standard deviations from the mean, respectively. The two-standard-deviation line is approximately aligned with the largest result from our human subjects.

also computed for each subject and each side separately and plotted together on a single graph. We assume the nominal value of K, $||\mathbf{s}||$, and θ for these results.

In Figs. 8, 9(b), and 10, we chose to present the 60 individual subject-side curves, rather than showing only aggregate statistics. This choice enables the reader to visualize the difference in r, as a function of the respective parameters, for the largest and smallest human subjects in our study, as well as the distribution generally.

7. Oversized MDS

We now consider the design of a one-size-fits-all MDS for magnetically guided insertions. In order to generate the necessary magnetic field across as many patients as needed, the MDS will be needed to be oversized for the vast majority of patients. As a result, a one-size-fits-all MDS will not need to get equally close to all patients' heads during the closest approach (i.e. at the point of maximum field generation).

To gain some insight into this, we choose a 75-mmradius MDS; this is only about 1.5 mm larger than the maximum value for r listed in Table A.3. All the valid locations generated by the brute-force method for each subject and each side are identified for the oversized MDS. Recall that valid locations are those in which two constraints must be satisfied. First, the MDS centered at this location must be able to generate the required magnetic field at the cochlea. Second, the MDS must not collide with any part of the subject's head at this location. Let us define the point cloud that contains all the valid points that are also anterior to the cochlea. We can compute the centroid of this point cloud, v, by averaging these valid points; the results are presented in Table A.5 of the appendix. Attempting to place the MDS at the centroid will yield a location that is robust (i.e. insensitive) to registration and other errors.

An example using P28 is provided in Fig. 11, where every point represents a valid location of the 75 mm oversized MDS in its closest approach to the subject's head. All points that are posterior to the cochlea are shown in red; all points anterior to the cochlea are



Fig. 11. All valid locations to place an MDS that has been oversized to 75 mm radius, for P28. Black points are anterior to the cochlea. Red points are posterior to the cochlea. The black dots inside the head (left) locate the cochleae.

shown in black. Since the optimal radius for P28 is approximately 59 mm, an MDS whose radius is 75 mm represents a 28% increase in the radius over what is needed. This yields many valid points at which the MDS can be located. In practice, it is best to ignore all the locations that are posterior to the cochlea because the insertion is conducted behind the ear, and placing an MDS anywhere behind the ear is impractical. Further, potential points superior and inferior to the top-most and bottom-most CT slice are ignored since there is no sufficient surface data available for collision detection.

In general, the shape of the point cloud also suggests more robustness along the y- and z-directions as compared to the x-direction. In principle, this should not be a problem because the initial alignment of the MDS requires the edge of the MDS to contact the surface of the head. This physical constraint will help enforce the tighter tolerance needed in the x-direction.

8. Electromagnetic MDS

Until this point, we have exclusively considered a spherical permanent magnet as the MDS, due to both the simplicity in modeling the field generated by a spherical permanent magnet, as well as the simplicity of its geometry for collision detection. However, it may be desirable to use an electromagnetic source as the MDS for three reasons: First, an electromagnet has a controllable magnetic dipole, which means that it does not need to be actively translated during the insertion to vary the field strength at the cochlea. Second, an electromagnet can be turned off, making it inert when not in use; a strong permanent magnet represents an ongoing safety concern. Third, the relatively short duration of an electrode-array insertion (typically less than 30 sec) would allow high levels of current to be sourced through the coils without reaching unsafe temperatures.

Our group previously developed an electromagnetic source that was explicitly optimized to be modeled as a magnetic point-dipole source: the Omnimagnet [12]. The Omnimagnet comprises three nested mutually orthogonal coils, with a spherical ferromagnetic core at the common center. The geometries of the individual coils and the size of the spherical core were parametrically optimized to fit within a cubic bounding volume of sidelength L (in units $\{m\}$) while achieving the following specifications: First, the dipole strength in each of the three orthogonal basis directions is the same given the same current density J (in units $\{A \cdot m^{-2}\}$) in each coil. Second, the second-largest term in the multipole expansion after the dipole field, which is called the quadrupole field, is zero in all three basis directions to ensure that the dipole-field model accurately approximates the Omnimagnet's field at close distances. The strength of the dipole field decays as $\|\mathbf{p}\|^{-3}$, the quadrupole field

decays as $\|\mathbf{p}\|^{-5}$, and the remaining terms in the multipole expansion decay as $\|\mathbf{p}\|^{-7}$, $\|\mathbf{p}\|^{-9}$, etc. Third, the Omnimagnet has as strong a dipole as possible for a given current density in the coils.

The field of an Omnimagnet can be approximated by the point-dipole model (3) outside of its minimum bounding sphere (which touches the corners of the cube and has a radius of $\sqrt{3}L$), and is accurate to less than 5% error at all points outside of 1.5 minimum-boundingsphere radii [12]. The dipole moment of an Omnimagnet is $\|\mathbf{M}\| = 0.051L^4 J$ [12]. In any actual engineered Omnimagnet, there will be limits on the current density *J* that can be sourced through the coils due to the power supply and amplifiers chosen as well as safe heating limits. At this maximum J, the Omnimagnet's dipole strength will be at its maximum, and we can use this maximum dipole strength to make a direct analogy to the sphericalpermanent-magnet dipole strength that we have used throughout the development in this paper. We can use whatever bounding-sphere radii we choose to achieve a given level of accuracy, which will change the effective values of \tilde{M} and B_r that are distributed over the effective spherical volume. This method will tend to be conservative, since the patient's head could have been allowed to partially reside inside of the bounding sphere while still avoiding collision with the Omnimagnet.

9. Discussion

In this paper, we have rigorously described the process by which anyone could determine, on their own, the minimum MDS size and its location for a given patient, and for a given EAT magnet strength and required torque. For a device maker interested in building an MDS to accommodate the general population, it might be worthwhile to obtain a larger number of head surfaces (particularly of large heads) to supplement the data set examined here. For this purpose, our data set of rendered surface points and surface normals is available upon request by contacting the corresponding author. Our description of a 75 mm MDS in Sec. 7 should not be construed as a recommendation of the optimal one-sizefits-all size. First, the dimension was arbitrarily chosen to slightly exceed the maximum r in Table A.3. It may be more desirable to specify a size from an average and standard deviation based on this study or a larger population study. Second, the results in Sec. 7 assumed no clearance between the head surface and the magnet (i.e. $\delta = 0$) and assumed the nominal value of K. In practice, the available magnetic grade (see Fig. 9) and the designed housing thickness (see Fig. 10) will drastically impact the size of a permanent-magnet MDS, and the allowable current density will drastically impact the size of an electromagnetic MDS. Third, we have determined in our prior work [8] that the required magnetic field can vary substantially based on electrode-array models. Therefore, it may even be necessary to have electrodearray specific MDS models as well. Even the progress toward thinner and more flexible electrode arrays may yield a smaller MDS in the future. Fortunately, these electrode-array specific variations, whether magnetic or mechanical, can be handled by determining their effect on the required magnetic field $||\mathbf{B}||$ (see Fig. 9). Finally, a specific registration tolerance might be preferred by the clinicians and would require an analysis similar to that done for Fig. 8. Under these various circumstances, the device maker can either conduct a full optimal configuration analysis (Sec. 4) or use the sensitivity curves in Sec. 6 for their MDS design.

The magnetic guidance strategy works best if both the lumen and modiolar axis are determined for each patient. Fortunately, it is standard practice to include a CT-based radiological assessment as part of the complete preoperative medical assessment for cochlear-implant candidacy. There is little reason for not planning patientspecific magnetically guided insertions given the ease in which cochlea segmentation can be done [16]. Further, since software can be implemented to generate a patientspecific head surface (as described in Sec. 3), the method outlined in Sec. 7 can be used to determine a patientspecific centroid location for a one-size-fits-all MDS that is most robust to registration error. Although our work describes this centroid with respect to the cochlea, it can be expressed relative to external markers on the patient (e.g. bone anchors [21]). In such a scenario, an optical tracker could register the MDS to the centroid location with high accuracy [22].

In traditional ear surgery, a large retractor is used to reflect the ear forward so that the surgeon can look inside the area with a microscope. It is possible that this could interfere with the MDS placement, particularly in the case of the spherical-permanent-magnet MDS; the cubic Omnimagnet can be rotated to mitigate interference. However, since the MDS is intended to be used with an automated electrode-insertion tool, the retractor could be removed since visualizing the facial recess is unnecessary to accomplish automated electrode insertions.

Standard temporal-bone CT protocols start the scan below the mastoid process and end just above the petrous ridge [23]. This effectively yields data sufficient only to render the middle portion of the head, typically including the eyes and nose. With the exception of P28, all CT scans of the subjects were done this way. Only the cadaver scans (C1 and C2) included the section below the nose to the shoulder. Fortunately, the optimal magnet size and location likely would not change with more slices. This can be seen in Fig. 12, where we show contour maps of distances from the cochlea to the surface points for C1 and C2, as well as for one standard-protocol



Fig. 12. Contour map of ||s|| for P25 and C1. Blue dots indicate where the optimal MDS touches the surface of the head.

subject (P25); in addition, the point where the optimal MDS touches the surface is indicated by a large blue dot. The smallest possible magnet will be located, where the distance is shortest to the cochlea, barring collision with the head, and it is clear that using additional CT data will not find a global optimal that is different than the local optimal found using the standard-protocol data.

10. Conclusion

We described a method to determine the patient-specific configuration (size and location) of a sphericalpermanent-magnet dipole-field source for magnetically guided cochlear-implant electrode-array insertions, using standard CT temporal bone scans. To generate 100 mT at the cochlea, the optimal configuration of the dipole-field source should be lateral-to and slightly anterior-to the cochlea with an approximate radius (mean and standard deviation across subjects) of 64 mm and 4.5 mm, respectively. When compared to the configuration assumed in prior work, the optimal location yields a 43% reduction in the dipole-field source's radius and a nearly fivefold reduction in its volume. Although potential magnetic forces that may direct the tip into the basilar membrane are increased by nearly 80% at the optimal configuration compared to the modiolar configuration, they are still two orders of magnitude below the threshold needed to puncture the membrane. Sensitivity curves were generated for the minimum radius at the optimal configuration and demonstrate that the optimization is most sensitive to the magnetic field requirement. A one-size-fits-all configuration with a dipole-field source of approximately 75 mm will yield a solution that is robust to registration error.

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Appendix A. Tabular Data

In this appendix, we provide tabulated information used and data obtained in this paper. Table A.1 provides the MATLAB commands used in the image-processing steps of Sec. 3. Table A.2 provides information related to the 30 human-subject CT scans used in Sec. 3. Table A.3 describes the optimal size and placement of the MDS for each human subject, as well as the aggregate statistics, obtained in Sec. 4.2.1. It also includes the results for the modiolar configuration, obtained in Sec. 4.2.2. Table A.4 shows analogous results using the brute-force approach of Sec. 4.2.3. Table A.5 shows the location of a 75-mmradius one-size-fits-all MDS, optimized for each human subject to be robust to registration errors, obtained in Sec. 7.

Table A.1. MATLAB commands used in the image-processing steps described in Sec. 3.

Image-Processing Step	MATLAB Command
Convert DICOM data Black-and-white conversion Skull boundary segmentation Skull boundary mask	dicomread and dicominfo im2bw and graythresh bwboundaries or bwtraceboudaries poly2mask
Filtering Surface interpolation Surface normals	imerode smooth3 isosurface isonormals

Table A.2.	Database of human-subject CT scans. P1–P28 are scans of anonymous
subjects pro	ovided by the University of Utah (UU). C1-C2 are scans of cadavers
obtained fro	om the National Library of Medicine (NLM).

	Gender	Age		Resoluti	Total		
ID	(M/F)	(yrs)	Source	Pixel	Slice	Slices	
P1	F	62	UU	0.41	0.6	89	
P2	F	52	UU	0.41	2	39	
Р3	F	33	UU	0.41	2	37	
P4	F	60	UU	0.41	2	37	
P5	М	26	UU	0.44	5	16	
P6	F	42	UU	0.41	2	30	
P7	F	85	UU	0.41	4	16	
P8	F	21	UU	0.39	5	15	
P9	F	32	UU	0.41	5	15	
P10	F	65	UU	0.47	0.6	115	
P11	F	39	UU	0.49	3	23	
P12	М	48	UU	0.44	2	39	
P13	F	51	UU	0.43	1	68	
P14	Μ	64	UU	0.46	5	17	
P15	F	47	UU	0.34	2	25	
P16	М	79	UU	0.59	2	36	
P17	F	48	UU	0.43	2	33	
P18	F	52	UU	0.33	0.7	82	
P19	F	29	UU	0.57	2	32	
P20	F	39	UU	0.41	2	31	
P21	F	37	UU	0.59	2	36	
P22	F	36	UU	0.38	5	12	
P23	F	60	UU	0.41	0.6	98	
P24	Μ	36	UU	0.46	2	31	
P25	F	67	UU	0.39	1	78	
P26	Μ	83	UU	0.41	1	47	
P27	F	44	UU	0.32	2	31	
P28	М	8	UU	0.45	1	161	
C1	Μ	—	NLM	0.79	1.45	183	
C2	F	—	NLM	0.54	1	249	

	r_m	102	113	119	105	105	105	103	115	117	121	116	107	114	113	99.1	117	100	116	118	121	105	115	117	121	110	131	113	124	114	109	99.1	131	31.9	113 7.58
	r	62.4	67.1	61.2	62.8	65.1	63.5	67.4	61.9	59.0	66.2	67.5	71.4	65.3	69.4	55.0	64.3	58.5	60.0	72.7	66.1	59.8	64.8	65.4	59.2	58.1	60.2	61.9	58.9	67.6	71.0	55.0	72.7	17.7	63.8 4.39
	θ	45.5	31.7	37.5	31.8	31.6	35.1	20.5	43.3	29.8	39.3	24.7	27.8	22.7	23.8	17.5	27.3	36.4	37.6	24.9	38.9	38.2	37.2	32.1	32.3	16.9	24.3	36.5	29.9	31.3	22.0	16.9	45.5	28.6	30.9 7.41
	$s_z]^T$	-0.23]	0.18]	0.05]	-0.08]	0.42]	-0.24]	-0.10]	-0.20]	-0.28]	0.10]	0.36]	-0.32]	-0.34]	-0.39]	-0.20]	-0.19]	-0.35]	-0.19]	0.22]	-0.18]	0.03]	-0.12]	-0.10]	0.04]	-0.18]	0.13]	-0.04]	-0.20]	0.25]	-0.31]				-0.08] 0.76]
de	sy	- 89.0 (0.58	0.46	0.51 -	0.14	0.61 -	3 0.19 -	99.0 8	ł 0.46 –	0.76	0.23	0.31 -	0.32 -	ł 0.37 –	5 0.17 -	7 0.61 -	0.53 -	0.55 -	3 0.42	7 0.61 -	3 0.63	0.71 -	0.48 -	0.59	0.24 -	0.43	0.64 -	3 0.43 -	5 0.44	2 0.25 -				7 0.49 – 3 0.59
Right Si	$[s_x$	[0.70	[0.8([0.89	[0.86	[0.85	[0.75	36.0]	[0.73	[0.87	[0.65	[0.9([0.89	[0.89	[0.84	[0.96	[0.77	[0.77	[0.8]	[0.8	[0.77	[0.78	<u>[0.65</u>	[0.87	[0.8(36.0]	[0.89	[0.77	[0.8	[0.86	[0.92				[0.87 [0.28
	s	52.3	51.1	48.3	47.9	49.6	49.4	48.8	51.0	44.5	53.0	49.7	53.3	47.7	50.9	39.4	47.9	45.9	47.4	53.5	52.7	47.5	51.1	50.0	45.3	41.6	44.2	48.6	44.5	51.4	51.6	39.4	53.5	14.2	48.7 3.49
	$\hat{p}_z]^T$	-0.22]	0.15]	-0.07	-0.08]	0.12]	-0.13]	0.10]	-0.10]	-0.15]	0.03]	0.15]	[60.0-	-0.13]	-0.17]	-0.08]	-0.10]	-0.19]	-0.10]	0.16]	-0.19]	0.01]	[60.0-	[0.0.0]	-0.02]	-0.01	0.02]	[0.0.0]	-0.17]	0.12]	-0.11]				-0.04] 0.73]
	$\int_{x} \hat{p}_{y}$	2 0.31	4 0.30	9 0.13	7 0.22	8 0.16	3 0.34	9 0.14	4 0.32	6 0.23	8 0.47	8 0.12	8 0.19	6 0.25	2 0.34	0 0.04	1 0.40	5 0.26	7 0.24	7 0.19	4 0.28	5 0.30	0 0.42	8 0.20	5 0.32	6 0.27	7 0.23	4 0.33	7 0.15	7 0.19	8 0.17				7 0.25 0 0.66
	-[<i>þ</i>	[0.9	[0.9	6.0]	[0.9]	[0.9]	[0.9]	[0.9	[0.9]	[0.9	[0.8	6.0]	[0.9	[0.9	[0.9]	[1.0	[0.9]	6.0]	[0.9]	6.0]	[0.9]	[0.9	[0.9]	6.0]	6.0]	6.0]	[0.9	[0.9	[0.9]	[0.9]	[0.9			-	[0.9 [0.2
	b	106	114	104	107	111	108	114	105	100	112	115	121	111	118	93.3	109	99.3	102	123	112	101	110	111	100	98.6	102	105	100	115	120	93.3	123	30.0	108 7.45
	r_m	109	94.1	127	97.6	125	131	114	115	103	118	113	107	109	109	98.0	132	110	114	109	105	97.3	113	115	112	96.5	128	106	115	114	114	94.1	132	37.6	$112 \\ 10.0$
	r	59.6	69.1	66.8	62.7	71.3	62.2	66.7	63.9	62.4	61.6	67.1	6.9.9	69.1	71.7	60.4	62.2	60.3	60.3	69.69	64.5	56.1	63.0	63.6	59.1	58.4	60.8	61.6	58.3	73.5	72.0	56.1	73.5	17.4	64.3 4.75
	θ	35.7	21.7	25.8	30.8	23.4	36.7	17.2	42.8	33.0	35.1	34.7	40.7	31.3	32.3	25.5	34.8	36.5	34.8	26.5	37.9	40.6	16.3	21.6	32.5	36.7	15.1	26.3	20.0	35.8	15.8	15.1	42.8	27.6	29.9 8.07
	$s_z]^T$	-0.22]	-0.21]	-0.28]	0.13]	-0.21]	-0.16]	-0.02]	-0.38]	-0.36	0.13]	0.29]	0.52]	-0.17]	[60.0]	-0.27]	-0.06]	-0.42]	-0.11]	0.01]	-0.17]	-0.16]	-0.30]	-0.32]	-0.04]	[60.0]	-0.27]	-0.19]	-0.19]	0.21]	0.04]				-0.10] 0.79]
	sy	3 0.59 -	0.56 -	3 0.25	2 0.56	5 0.22	4 0.65 -	4 0.34 -	2 0.58 -	5 0.39 -	5 0.65	4 0.46	3 0.45	1 0.56 -	3 0.68	0.34	3 0.68	5 0.49 -	2 0.57 -	2 0.39	7 0.61 -	- 09.0 6	1 0.28 -	3 0.19 -	1 0.58 -	1 0.70	4 0.20	3 0.43 -	0.40	4 0.51	3 0.37				5 0.50 0.54
eft Sid€	$[s_x$	[-0.78	[-0.8([-0.9]	[-0.8]	[-0.9]	[-0.7]	-0.9	[-0.7]	[-0.8]	[-0.7]	[-0.8	[-0.7]	[-0.8]	[-0.7]	[-0.9([-0.7]	[-0.76	[-0.8]	[-0.9]	[-0.7]	[-0.79]	[-0.9]	[-0.9]	[-0.8]	[-0.7]	-0.9	[-0.8]	[-0.9([-0.8]	[-0.9]				[-0.86 [0.29
Γ	s	46.5	50.2	49.4	47.6	52.1	48.9	47.7	52.4	47.9	47.9	52.1	56.4	52.5	54.9	44.6	48.3	47.3	46.8	51.7	51.2	45.3	44.9	46.2	45.3	45.9	43.2	45.7	42.1	57.4	51.3	42.1	57.4	15.3	48.8 3.78
	$\hat{p}_{z}]^{T}$	-0.10]	-0.15]	-0.06]	0.06]	-0.02]	-0.12]	0.11]	-0.21]	-0.12]	0.11]	0.13]	0.21]	-0.07]	0.05]	-0.08]	-0.06]	-0.25]	-0.06]	0.15]	-0.21]	-0.18]	-0.15]	-0.11]	-0.02]	0.04]	-0.15]	0.04]	-0.17]	0.08]	0.02]				-0.04] 0.79]
	$\hat{p_y}$	0.31 -	0.38 -	0.13 -	0.30	- 60.0	0.35 -	0.22	0.28 -	0.21 -	0.36	0.19	0.31	0.31 -	0.41	0.18 -	0.40 -	0.24 -	0.26 -	0.17	0.28 -	0.24 -	0.24 -	0.16 -	0.29 -	0.41	0.11 -	0.32	0.21 -	0.21	0.21				0.26 - 0.59
	$-[\hat{p_x}]$	[-0.95]	[-0.91	[-0.99	[-0.95]	[-1.00]	[-0.93	[-0.97]	[-0.94	[-0.97]	[-0.93	[-0.97]	[-0.93	[-0.95]	[-0.91]	[-0.98	[-0.92]	[-0.94	[-0.96	[-0.97]	[-0.94	[-0.95]	[-0.96	[-0.98	[-0.96	[-0.91	[-0.98	[-0.95]	[-0.96	[-0.97]	[-0.98				[-0.96]
	b	101	117	113	106	121	106	113	108	106	105	114	119	117	122	102	106	102	102	118	110	95.1	107	108	100	99.0	103	105	98.9	125	122	95.1	125	29.6	109 8.05
	ID	P1	P2	P3	P4	P5	P6	P7	P8	P9	P10	P11	P12	P13	P14	P15	P16	P17	P18	P19	P20	P21	P22	P23	P24	P25	P26	P27	P28	C1	C2	Min	Max	Range	Avg Std

Table A.4. Optimal configuration of MDS as defined by its radius (*r*) and its location relative to the cochlea (**p**), using the brute-force method. Vectors are expressed in the RAS convention, where +*x* is directed to the person's right side, +*y* is directed anteriorly, and +*z* is directed superiorly. Units are in {mm} except for the unitless direction vector $-\hat{\mathbf{p}} = -[\hat{p}_x \quad \hat{p}_y \quad \hat{p}_z]^T$ from the cochlea toward the MDS center.

		Left		Right							
ID	$\ \mathbf{p}\ $	$-[\hat{p}_x \hat{p}_y \hat{p}_z]^T$	r	 p 	$-[\hat{p}_x \hat{p}_y \hat{p}_z]^T$	r					
P1	100	[-0.96 0.28 -0.08]	58.8	100	[0.95 0.28 -0.13]	59.1					
P2	112	[-0.95 0.32 0.00]	65.8	114	[0.94 0.30 0.14]	67.1					
Р3	112	[-0.99 0.12 -0.04]	66.1	103	[0.99 0.12 -0.09]	60.8					
P4	105	[-0.96 0.28 0.04]	62.1	106	$[0.98 \ 0.20 \ -0.09]$	62.7					
P5	120	[-0.99 0.11 0.04]	70.9	109	[0.99 0.13 0.09]	64.3					
P6	103	$[-0.94 \ 0.29 \ -0.17]$	60.9	105	$[0.96 \ 0.23 \ -0.17]$	61.6					
P7	113	[-0.98 0.17 0.14]	66.4	112	[0.97 0.14 0.18]	66.1					
P8	105	$[-0.93 \ 0.16 \ -0.33]$	62.1	104	$[0.93 \ 0.17 \ -0.34]$	61.5					
P9	105	$[-0.98 \ 0.18 \ -0.12]$	62.0	100	$[0.96 \ 0.22 \ -0.17]$	59.0					
P10	104	[-0.94 0.34 0.08]	61.4	100	[0.95 0.29 0.06]	58.8					
P11	113	[-0.98 0.15 0.11]	66.8	113	[0.99 0.07 0.08]	66.7					
P12	111	[-0.96 0.26 0.07]	65.6	114	[0.96 0.26 0.07]	67.1					
P13	117	$[-0.96 \ 0.27 \ -0.11]$	68.8	109	[0.96 0.26 -0.09]	64.5					
P14	117	$[-0.94 \ 0.31 \ -0.13]$	69.0	117	$[0.93 \ 0.34 \ -0.13]$	68.9					
P15	102	$[-0.98 \ 0.05 \ -0.18]$	60.4	93.1	$[1.00 \ 0.08 \ -0.02]$	54.9					
P16	104	$[-0.93 \ 0.37 \ -0.06]$	61.1	110	$[0.92 \ 0.38 \ -0.10]$	64.6					
P17	101	$[-0.95 \ 0.21 \ -0.22]$	59.4	97.4	[0.96 0.21 -0.21]	57.4					
P18	101	[-0.96 0.23 -0.13]	59.7	101	[0.97 0.19 -0.16]	59.5					
P19	118	[-0.99 0.14 0.08]	69.8	123	[0.98 0.18 0.13]	72.2					
P20	109	$[-0.94 \ 0.26 \ -0.20]$	64.0	110	[0.96 0.24 -0.15]	64.8					
P21	94.9	$[-0.96 \ 0.22 \ -0.19]$	55.9	97.2	[0.95 0.21 -0.23]	57.3					
P22	105	[-0.95 0.32 0.01]	62.1	102	[0.94 0.31 -0.12]	60.2					
P23	103	$[-0.99 \ 0.14 \ -0.01]$	61.0	108	[0.99 0.14 -0.03]	63.9					
P24	100	[-0.96 0.28 -0.02]	59.1	101	[0.95 0.31 -0.04]	59.3					
P25	94.4	$[-0.94 \ 0.35 \ -0.03]$	55.7	97.5	[0.95 0.31 0.02]	57.5					
P26	102	$[-0.98 \ 0.18 \ -0.07]$	60.4	101	[0.98 0.21 0.01]	59.6					
P27	104	[-0.95 0.30 0.06]	61.0	101	[0.96 0.28 0.08]	59.4					
P28	97.9	$[-0.97 \ 0.16 \ -0.16]$	57.7	99.2	[0.98 0.13 -0.16]	58.5					
C1	123	[-0.98 0.20 0.03]	72.7	114	[0.99 0.15 0.07]	67.5					
C2	122	$[-0.98 \ 0.18 \ -0.02]$	71.9	119	[0.99 0.10 -0.11]	70.2					
Min	94.4		55.7	93.1		54.9					
Max	123		72.7	123		72.2					
Rng	28.9		17.1	29.4		17.4					
Avg	107	$[-0.97 \ 0.23 \ -0.05]$	63.3	106	$[0.98 \ 0.21 \ -0.05]$	62.5					
Std	7.89	[0.15 0.58 0.80]	4.65	7.40	[0.13 0.52 0.84]	4.36					

Table A.5. Placement, relative to the cochlea, of an MDS oversized to 75 mm radius. Vectors are expressed in the RAS convention. The values represent the centroid of all the valid points anterior to the cochlea, where the oversized MDS can be placed without collision and satisfy the magnetic field requirement. Units are in mm except for the unitless direction vector $\hat{\mathbf{v}} = [\hat{\mathbf{v}}_x \ \hat{\mathbf{v}}_y \ \hat{\mathbf{v}}_z]^T$.

		Left			Righ		
ID	$\ \mathbf{v}\ $	$[\mathbf{v}_x \mathbf{v}_y]$	$\mathbf{v}_{z}]^{T}$	$\ \mathbf{v}\ $	$[v_x]$	\mathbf{v}_y	$\mathbf{v}_{z}]^{T}$
P1	122	[-0.95 0.32	-0.03]	122	[0.96	0.29	-0.04]
P2	124	[-0.94 0.33	-0.01]	124	[0.95	0.29	0.08]
P3	125	[-0.99 0.16	-0.03]	123	[0.98	0.17	-0.08]
P4	123	[-0.96 0.29	0.01]	123	[0.97	0.24	-0.06]
P5	126	$[-0.99 \ 0.12$	0.04]	124	[0.98	0.20	0.05]
P6	122	[-0.94 0.31	-0.11]	122	[0.95	0.30	-0.08]
P7	124	$[-0.98 \ 0.20$	0.03]	124	[0.98	0.20	0.05]
P8	122	$[-0.97 \ 0.24$	-0.08]	121	[0.97	0.23	-0.05]
P9	122	$[-0.98 \ 0.21$	-0.03]	121	[0.97	0.24	-0.05]
P10	122	$[-0.92 \ 0.39$	0.01]	121	[0.95	0.32	0.02]
P11	125	$[-0.98 \ 0.15$	0.09]	125	[0.99	0.13	0.11]
P12	124	[-0.96 0.26	0.04]	125	[0.97	0.24	0.03]
P13	125	$[-0.96 \ 0.27$	-0.09]	123	[0.95	0.31	-0.09]
P14	125	$[-0.94 \ 0.34$	-0.08]	125	[0.93	0.36	-0.10]
P15	122	$[-0.97 \ 0.21$	-0.08]	120	[0.97	0.22	-0.07]
P16	123	$[-0.92 \ 0.39$	-0.03]	124	[0.91	0.40	-0.10]
P17	122	[-0.96 0.26	-0.10]	121	[0.96	0.28	-0.08]
P18	121	$[-0.97 \ 0.23$	-0.06]	121	[0.98	0.20	-0.07]
P19	125	$[-0.98 \ 0.16$	0.11]	126	[0.98	0.17	0.11]
P20	124	$[-0.95 \ 0.28$	-0.13]	124	[0.96	0.24	-0.13]
P21	121	$[-0.97 \ 0.25$	-0.02]	121	[0.96	0.28	-0.01]
P22	122	$[-0.95 \ 0.30$	-0.05]	122	[0.95	0.32	-0.05]
P23	123	$[-0.97 \ 0.23$	-0.05]	124	[0.98	0.20	-0.04]
P24	121	[-0.95 0.31	-0.03]	121	[0.94	0.33	-0.04]
P25	120	[-0.93 0.35	-0.04]	121	[0.95	0.31	-0.02]
P26	122	$[-0.98 \ 0.20$	-0.05]	121	[0.97	0.23	-0.02]
P27	122	$[-0.96 \ 0.27$	-0.01]	121	[0.96	0.27	0.00]
P28	120	$[-0.97 \ 0.24$	-0.07]	121	[0.98	0.20	-0.10]
C1	127	$[-0.98 \ 0.17$	0.04]	125	[0.98	0.15	0.11]
C2	126	[-0.98 0.19	0.01]	125	[0.98	0.18	-0.05]
Min	120			120			
Max	127			126			
Rng	6.44			6.04			
Avg	123	[-0.96 0.26	-0.03]	123	[0.97	0.25	-0.03]
Std	1.74	[0.21 0.74	0.64]	1.68	[0.21	0.64	0.74]

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